Ordinary and partial differential equations

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ODEs and PDEs

Time-evolving phenomena: IVP (Initial Value Problem), usually Ordinary Differential Equations

Space-constraint phenomena: BVP (Boundary Value Problem), usually Partial Differential Equations



Ordinary Differential Equations



Simple Initial Value Problem

Newtons' Second Law: F = ma

Acceleration depends linearly on the force exerted on the mass: $a = \frac{d}{dt}u = \frac{d^2}{dt^2}x = \frac{F}{m}$

Write velocity as vector \rightarrow first order derivative in time: $u(t) = (x(t), x'(t))^t$

u implicitly depends on time, easier to solve: $u'(t) = f(u(t)), u(0) = u_0$



Numerical treatment of differential equations

Initial value problem: $u'(t) = f(t, u(t)), \qquad u(0) = u_0, \qquad t > 0$

Boundary value problem:

$$u''(x) = f(x), \ x(0) = x_0, \ x(1) = x_1, \qquad x \in [0, 1]$$

General assumption: *f* has higher derivatives.

IVP stability: solutions corresponding to different u_0 values converge as $t \to \infty$. Criterium:

$$\frac{\partial}{\partial u}f(t,u) = \begin{cases} > 0 & unstable \\ = 0 & neutrallystable \\ < 0 & stable \end{cases}$$

Simple example: $f(t, u) = -\lambda u$, then $u(t) = u_0 e^{-\lambda t}$; stable if $\lambda > 0$



Finite difference approximation

We turn the continuous problem into a discrete one, by looking at finite time/space steps.

Assume all functions are sufficiently smooth, and use Taylor series:

$$u(t + \Delta t) = u(t) + u'(t)\Delta t + u''(t)\frac{\Delta t^2}{2!} + u'''(t)\frac{\Delta t^3}{3!} + \cdots$$

This gives for u':

$$u'(t) = rac{u(t+\Delta t)-u(t)}{\Delta t} + O(\Delta t^2)$$

So we approximate

$$u'(t) pprox rac{u(t+\Delta t)-u(t)}{\Delta t}$$

and the "truncation error" is $O(\Delta t^2)$.



Finite differences 2

How does this help? In u' = f(t, u) substitute

$$u'(t)
ightarrow rac{u(t+\Delta t)-u(t)}{\Delta t}$$

giving

$$\frac{u(t+\Delta t)-u(t)}{\Delta t}=f(t,u(t))$$

or

$$u(t + \Delta t) = u(t) + \Delta t f(t, u(t))$$

Let
$$t_0 = 0$$
, $t_{k+1} = t_k + \Delta t = \cdots = (k+1)\Delta t$, $u(t_k) = u_k$:

$$u_{k+1} = u_k + \Delta t f(t_k, u_k)$$

Discretization 'Explicit Euler' or 'Euler forward'.

Does this compute something close to the true solution? 'Discretization error'



Some error analysis

Local Truncation Error: assume computed solution is exact at step k, how wrong will we be at step k + 1?

$$u(t_{k+1}) = u(t_k) + u'(t_k)\Delta t + u''(t_k)\frac{\Delta t^2}{2!} + \cdots$$

= $u(t_k) + f(t_k, u(t_k))\Delta t + u''(t_k)\frac{\Delta t^2}{2!} + \cdots$
 $u_{k+1} = u_k + f(t_k, u_k)\Delta t$

So

$$L_{k+1} = u_{k+1} - u(t_{k+1})$$

= $u_k - u(t_k) + f(t_k, u_k) - f(t_k, u(t_k)) - u''(t_k) \frac{\Delta t^2}{2!} + \cdots$
= $-u''(t_k) \frac{\Delta t^2}{2!} + \cdots$

Global error is first order method: $E_k \approx \Sigma_k L_k \approx k \left(\Delta t^2/2!\right) \approx (T - t_0/\Delta t) \left(\Delta t^2/2!\right) \approx O(\Delta t)$



An Euler forward example

Consider $f(t, u) = -\lambda u$, exact solution $u(t) = u_0 e^{-\lambda t}$; stable if $\lambda > 0$

Explicit Euler scheme

$$u_{k+1} = u_k - \Delta t \lambda u_k = (1 - \lambda \Delta t) u_k = (1 - \lambda \Delta t)^k u_0$$

Then

$$\begin{array}{l} u_k \to 0 \text{ as } k \to \infty \\ \Leftrightarrow \quad |1 - \lambda \Delta t| < 1 \\ \Leftrightarrow \quad -1 < 1 - \lambda \Delta t < 1 \\ \Leftrightarrow \quad -2 < -\lambda \Delta t < 0 \\ \Leftrightarrow \quad 0 < \lambda \Delta t < 2 \\ \Leftrightarrow \quad \Delta t < 2/\lambda \end{array}$$

Conditionally stable



Implicit Euler

Or 'Euler backward':

$$u(t-\Delta t)=u(t)-u'(t)\Delta t+u''(t)rac{\Delta t^2}{2!}+\cdots$$

SO

$$u'(t) = rac{u(t) - u(t - \Delta t)}{\Delta t} + u''(t)\Delta t/2 + \cdots$$

Compute u'(t) = f(t, u(t)) as

$$\frac{u(t) - u(t - \Delta t)}{\Delta t} = f(t, u(t))$$

$$\Rightarrow \quad u(t) = u(t - \Delta t) + \Delta t f(t, u(t))$$

$$\Rightarrow \quad u_{k+1} = u_k + \Delta t f(t_{k+1}, u_{k+1})$$

Implicit equation for u_{k+1} !

Let
$$f(t, u) = -u^3$$
, then $u_{k+1} = u_k - \Delta t u_{k+1}^3$
needs nonlinear solver



Stability of Implicit Euler

Again the $f(t, u) = -\lambda u$ example:

$$u_{k+1} = u_k - \lambda \Delta t u_{k+1}$$

(1 + $\lambda \Delta t$) $u_{k+1} = u_k$
 $u_{k+1} = \left(\frac{1}{1+\lambda\Delta t}\right) u_k = \left(\frac{1}{1+\lambda\Delta t}\right)^k u_0$

If $\lambda > 0$ (stable equation), then $u_k \rightarrow 0$ for all values of λ and Δt : unconditionally stable.

Pro: larger time steps possible, no worries Con: implicit equation needs to be solved



Higher order methods

Runge-Kutta Adams-Bashforth Crank-Nicholson



Boundary value problems

Consider
$$u''(x) = f(x, u, u')$$
 for $x \in [a, b]$ where $u(a) = u_a$, $u(b) = u_b$ in 1D and

$$-u_{xx}(\bar{x}) - u_{yy}(\bar{x}) = f(\bar{x}) \text{ for } x \in \Omega = [0, 1]^2 \text{ with } u(\bar{x}) = u_0 \text{ on } \delta\Omega.$$
(1)

in 2D.



Approximation of 2nd order derivatives

Taylor series (write h for δx):

$$u(x+h) = u(x) + u'(x)h + u''(x)\frac{h^2}{2!} + u'''(x)\frac{h^3}{3!} + u^{(4)}(x)\frac{h^4}{4!} + u^{(5)}(x)\frac{h^5}{5!} + \cdots$$

and

$$u(x-h) = u(x) - u'(x)h + u''(x)\frac{h^2}{2!} - u'''(x)\frac{h^3}{3!} + u^{(4)}(x)\frac{h^4}{4!} - u^{(5)}(x)\frac{h^5}{5!} + \cdots$$

Subtract:

$$u(x + h) + u(x - h) = 2u(x) + u''(x)h^{2} + u^{(4)}(x)\frac{h^{4}}{12} + \cdots$$

so

$$u''(x) = \frac{u(x+h) - 2u(x) + u(x-h)}{h^2} - u^{(4)}(x)\frac{h^4}{12} + \cdots$$

Numerical scheme:

$$-\frac{u(x+h)-2u(x)+u(x-h)}{h^2}=f(x,u(x),u'(x))$$

(2nd order PDEs are very common!)



This leads to linear algebra

$$-\frac{u(x+h)-2u(x)+u(x-h)}{h^2}=f(x,u(x),u'(x))$$

Equally spaced points on [0, 1]: $x_k = kh$ where h = 1/n, then

$$-u_{k+1} + 2u_k - u_{k-1} = h^2 f(x_k, u_k, u_k')$$
 for $k = 1, ..., n-1$

Written as matrix equation:

$$\begin{pmatrix} 2 & -1 & & \\ -1 & 2 & -1 & \\ & \ddots & \ddots & \ddots \end{pmatrix} \begin{pmatrix} u_1 \\ u_2 \\ \vdots \end{pmatrix} = \begin{pmatrix} h^2 f_1 + u_0 \\ h^2 f_2 \\ \vdots \end{pmatrix}$$



Matrix properties

- Very sparse, banded
- Symmetric (only because 2nd order problem)
- Sign pattern: positive diagonal, nonpositive off-diagonal (true for many second order methods)
- Positive definite (just like the continuous problem)



Initial Boundary value problem

Heat conduction in a rod T(x, t) for $x \in [a, b]$, t > 0:

$$\frac{\partial}{\partial t}T(x,t) - \alpha \frac{\partial^2}{\partial x^2}T(x,t) = q(x,t)$$

- Initial condition: $T(x,0) = T_0(x)$
- Boundary conditions: $T(a, t) = T_a(t)$, $T(b, t) = T_b(t)$
- Material property: $\alpha > 0$ is thermal diffusivity
- Forcing function: q(x, t) is externally applied heating.

The equation u''(x) = f above is the steady state.



Discretization

Space discretization: $x_0 = a$, $x_n = b$, $x_{j+1} = x_j + \Delta x$ Time discretiation: $t_0 = 0$, $t_{k+1} = t_k + \Delta t$ Let T_j^k approximate $T(x_j, t_k)$

Space:

$$\frac{\partial}{\partial t}T(x_j,t) - \alpha \frac{T(x_{j-1},t) - 2T(x_j,t) + T(x_{j+1},t)}{\Delta x^2} = q(x_j,t)$$

Explicit time stepping:

$$\frac{T_{j}^{k+1} - T_{j}^{k}}{\Delta t} - \alpha \frac{T_{j-1}^{k} - 2T_{j}^{k} + T_{j+1}^{k}}{\Delta x^{2}} = q_{j}^{k}$$

Implicit time stepping:

$$\frac{T_j^{k+1} - T_j^k}{\Delta t} - \alpha \frac{T_{j-1}^{k+1} - 2T_j^{k+1} + T_{j+1}^{k+1}}{\Delta x^2} = q_j^{k+1}$$



Computational form: explicit

$$T_j^{k+1} = T_j^k + \frac{\alpha \Delta t}{\Delta x^2} (T_{j-1}^k - 2T_j^k + T_{j+1}^k) + \Delta t q_j^k$$

This has an explicit form:

$$\underline{\mathcal{T}}^{k+1} = \left(I + \frac{\alpha \Delta t}{\Delta x^2}\right) \underline{\mathcal{T}}^k + \Delta t \underline{q}^k$$



Computational form: implicit

$$T_j^{k+1} - \frac{\alpha \Delta t}{\Delta x^2} (T_{j-1}^k - 2T_j^k + T_{j+1}^k) = T_j^k + \Delta t q_j^k$$

This has an implicit form:

$$\left(I - \frac{\alpha \Delta t}{\Delta x^2} \kappa\right) \mathcal{I}^{k+1} = \mathcal{I}^k + \Delta t \mathcal{q}^k$$

Needs to solve a linear system in every time step



Stability of explicit scheme

Let $q\equiv$ 0; assume $T_{j}^{k}=eta^{k}e^{i\ell x_{j}}$; for stability we require |eta|< 1:

$$T_{j}^{k+1} = T_{j}^{k} + \frac{\alpha \Delta t}{\Delta x^{2}} (T_{j-1}^{k} - 2T_{j}^{k} + T_{j+1}^{k})$$

$$\Rightarrow \beta^{k+1} e^{i\ell x_{j}} = \beta^{k} e^{i\ell x_{j}} + \frac{\alpha \Delta t}{\Delta x^{2}} (\beta^{k} e^{i\ell x_{j-1}} - 2\beta^{k} e^{i\ell x_{j}} + \beta^{k} e^{i\ell x_{j+1}})$$

$$\Rightarrow \beta = 1 + 2 \frac{\alpha \Delta t}{\Delta x^{2}} [\frac{1}{2} (e^{i\ell \Delta x} + e^{-\ell \Delta x}) - 1]$$

$$= 1 + 2 \frac{\alpha \Delta t}{\Delta x^{2}} (\cos(\ell \Delta x) - 1)$$



$$\frac{\beta^{k+1}}{\beta^k} = 1 + 2\frac{\alpha \Delta t}{\Delta x^2} (\cos(\ell \Delta x) - 1)$$

To get $|\beta| < 1$:

•
$$2\frac{\alpha\Delta t}{\Delta x^2}(\cos(\ell\Delta x) - 1) < 0$$
: automatic

• $2\frac{\alpha\Delta t}{\Delta x^2}(\cos(\ell\Delta x) - 1) > -2$: needs $2\frac{\alpha\Delta t}{\Delta x^2} < 1$, that is

$$\Delta t < \frac{\Delta x^2}{2\alpha}$$

big restriction on size of time steps



Stability of implicit scheme

$$T_{j}^{k+1} - \frac{\alpha \Delta t}{\Delta x^{2}} (T_{j_{1}}^{k+1} - 2T_{j}^{k+1} + T_{j+1}^{k+1}) = T_{j}^{k}$$

$$\Rightarrow \beta^{k+1} e^{i\ell\Delta x} - \frac{\alpha \Delta t}{\Delta x^{2}} (\beta^{k+1} e^{i\ell x_{j-1}} - 2\beta^{k+1} e^{i\ell x_{j}} + \beta^{k+1} e^{i\ell x_{j+1}}) = \beta^{k} e^{i\ell x_{j}}$$

$$\Rightarrow \beta^{-1} = 1 + 2\frac{\alpha \Delta t}{\Delta x^2} (1 - \cos(\ell \Delta x))$$
$$\beta = \frac{1}{1 + 2\frac{\alpha \Delta t}{\Delta x^2} (1 - \cos(\ell \Delta x))}$$

Noting that $1 - \cos(\ell \Delta x) > 0$, the condition $|\beta| < 1$ always satisfied: method always stable.



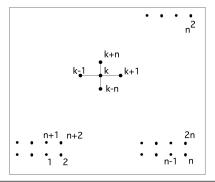
Sparse matrix in 2D case

Sparse matrices so far were tridiagonal: only in 1D case.

Two-dimensional: $-u_{xx} - u_{yy} = f$ on unit square $[0, 1]^2$

Difference equation:

 $4u(x,y) - u(x+h,y) - u(x-h,y) - u(x,y+h) - u(x,y-h) = h^2 f(x,y)$





Sparse matrix from 2D equation

